

# Spin-orbit torques

in ferromagnets and antiferromagnets

K Výborný, L P Zârbo, E K Vehstedt, J Sinova, V Novák, T Jungwirth

*Institute of Physics, Academy of Sciences of the Czech Rep., Prague*

*Dept. of Physics, Texas A&M University, College Station*

*Inst. f. Physik, Johannes Gutenberg Univ. Mainz*

H Li, A Manchon

*Physical Science and Engineering Division, KAUST*

F Freimuth

*Inst. f. Advanced Simulation, Forschungszentrum Juelich*

P Wadley, R P Campion, B L Gallagher

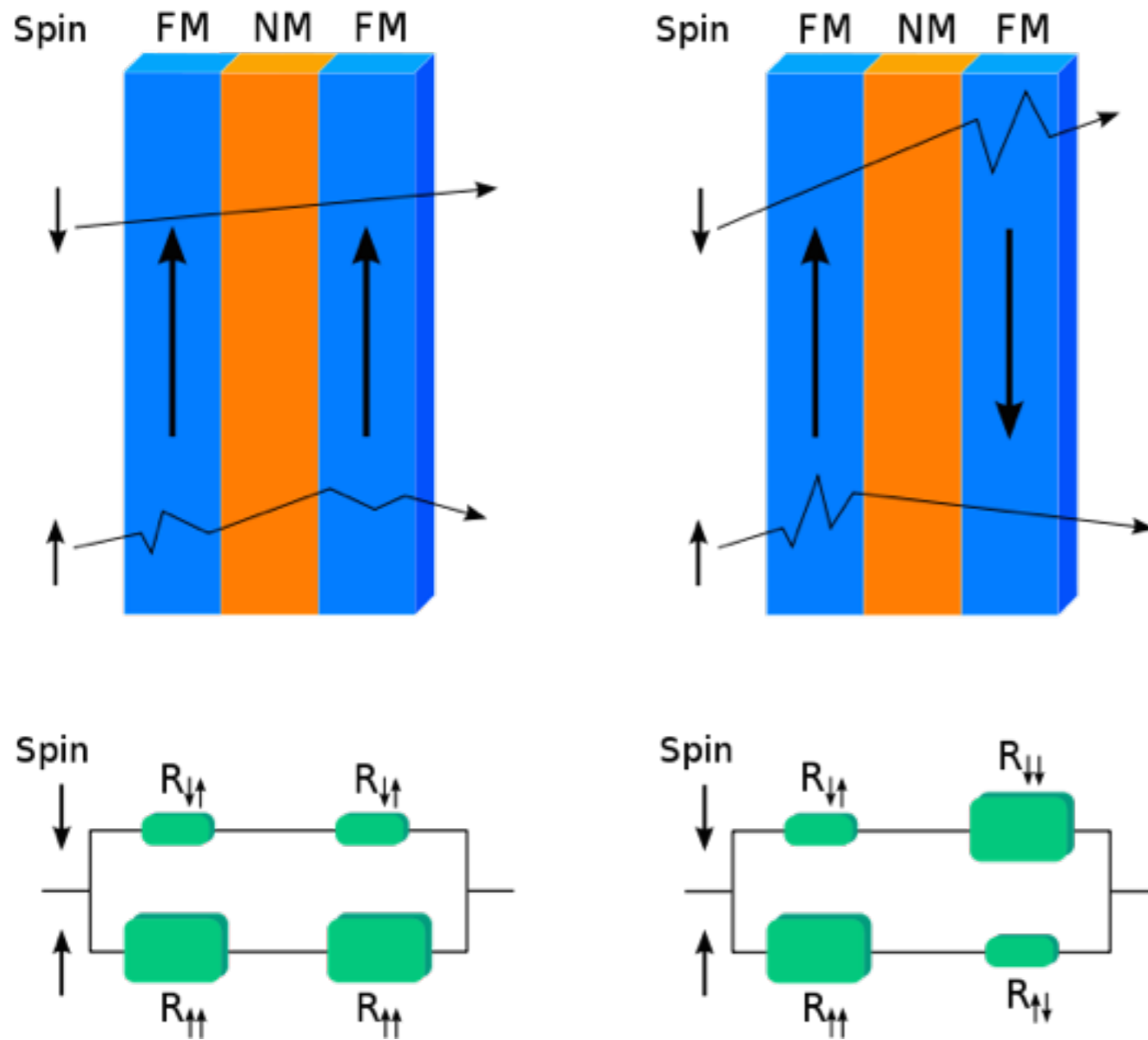
*School of Physics and Astronomy, University of Nottingham*

H Kurebayashi, D Fang, A C Irvine, T D Skinner, J Wunderlich, A J Ferguson

*Cavendish Laboratory, University of Cambridge*

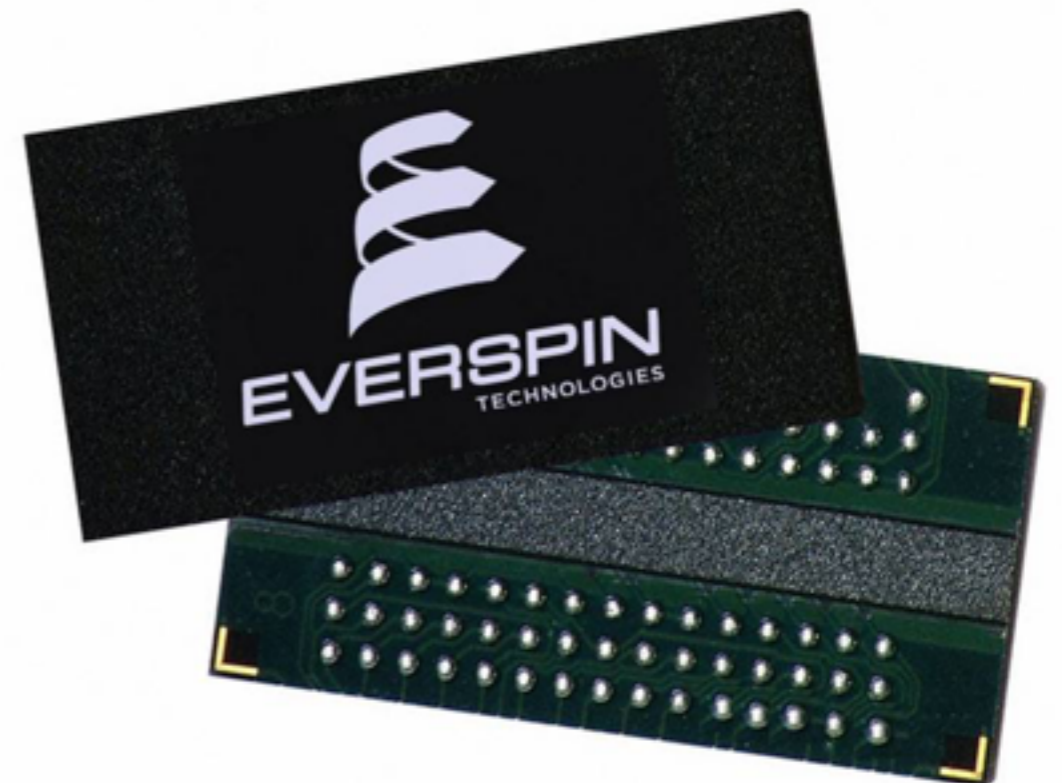
# Memory device (one possible concept)

## Giant Magnetoresistance (GMR)



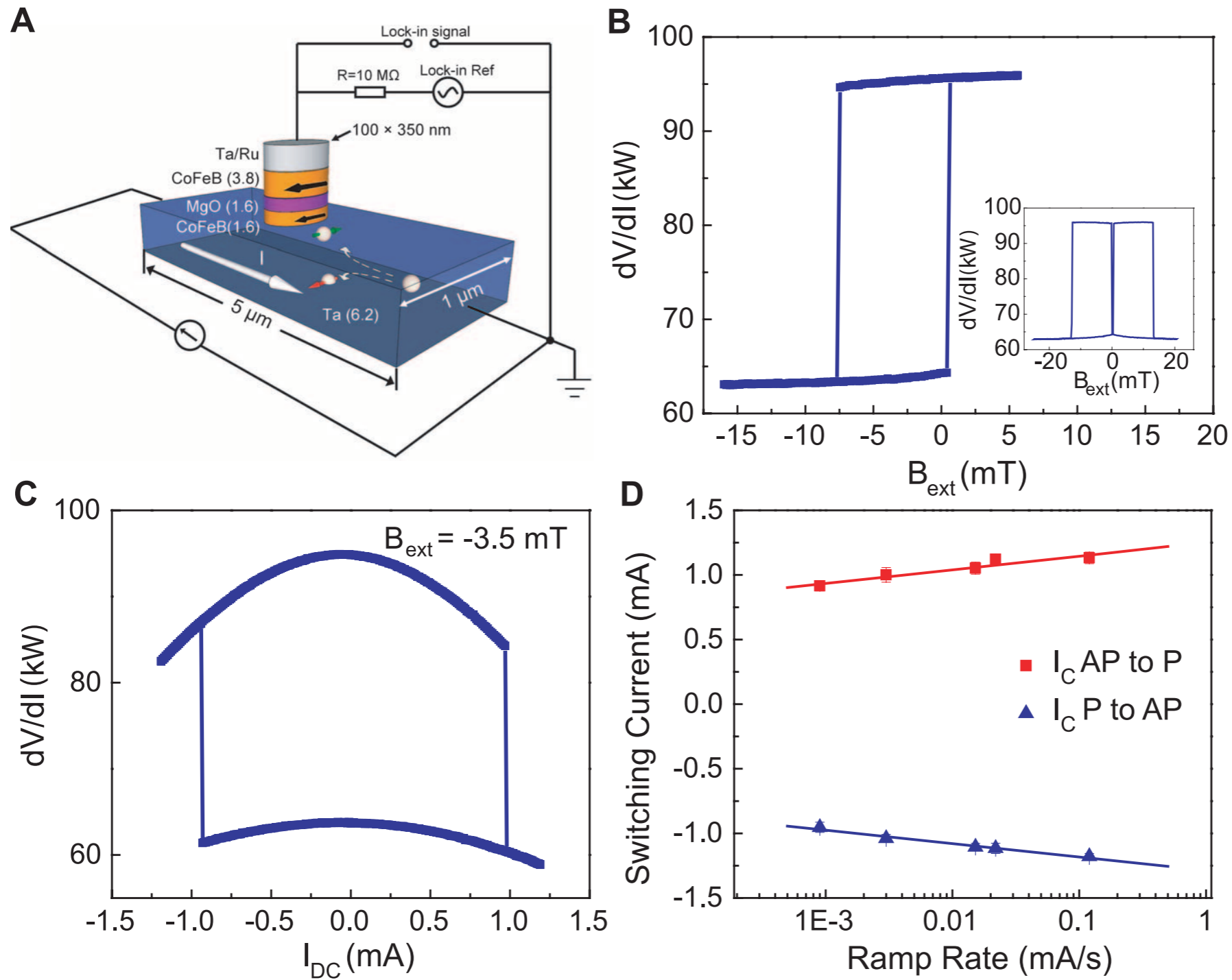
alternative:

- magnetic tunnel junctions
- how to write?
- spin transfer torque



[wikipedia]

# New developments - switching by SHE-spin torque



Ralph&Buhrman, Science 336, 555 (2012)

# Magnetic tunnel junctions as memory elements

writing information:

- spin-transfer torque
- using spin Hall effect (SHE-torque)
- spin-orbit torque

reading information:

- tunnelling magnetoresistance (similar to GMR)
- anisotropic magnetoresistance (AMR)
- ... and more (tunnelling AMR)

# Magnetic tunnel junctions as memory elements

## writing information:

- spin-transfer torque
- using spin Hall effect (SHE-torque)
- spin-orbit torque



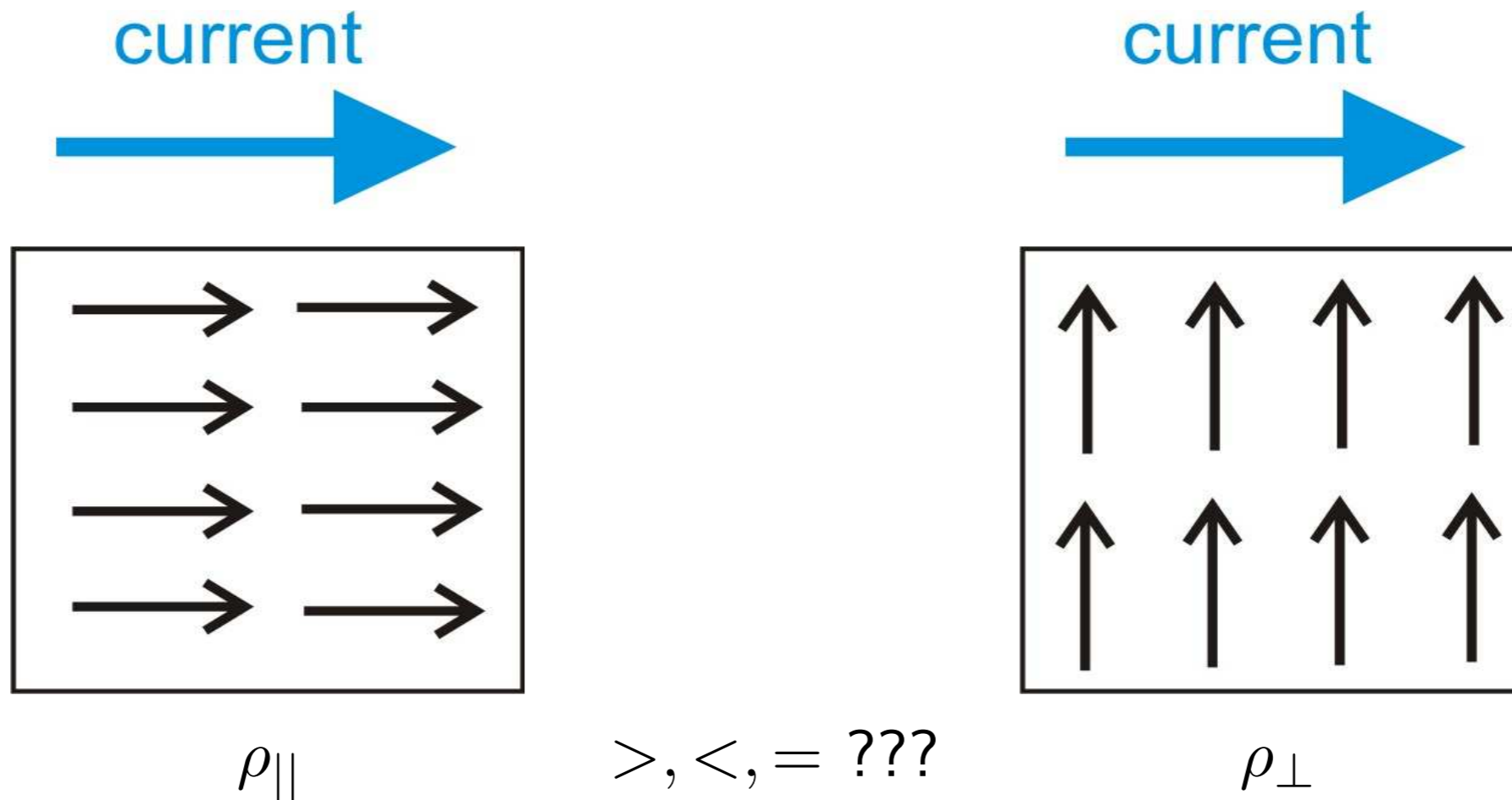
current-induced 'effective mag. field'

## reading information:

- tunnelling magnetoresistance (similar to GMR)
- anisotropic magnetoresistance (AMR)
- ... and more (tunnelling AMR)

# Anisotropic magnetoresistance (AMR)

the resistance depends on magnetization direction



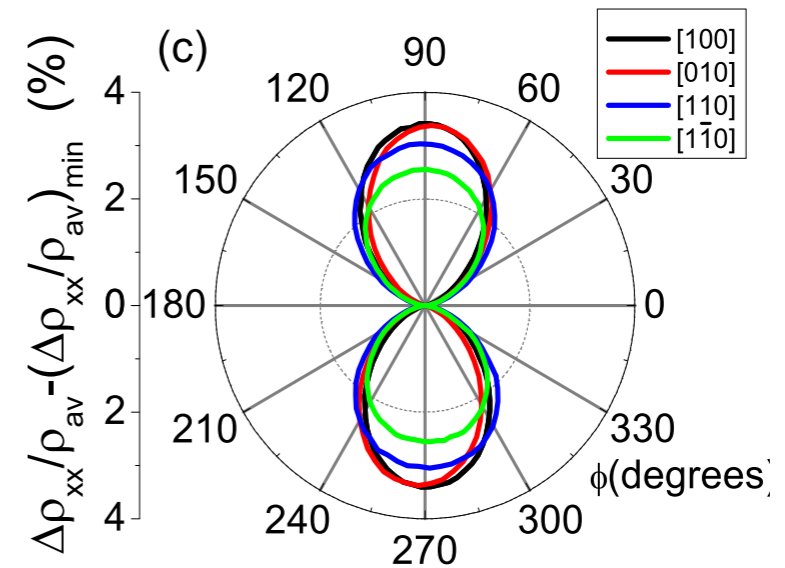
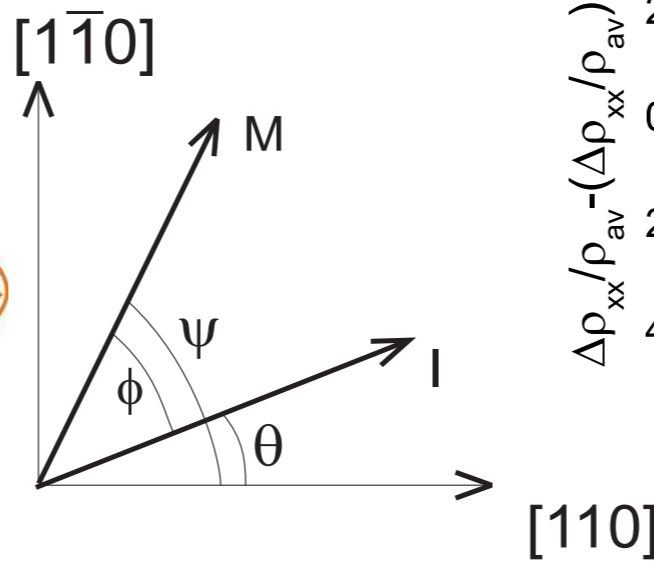
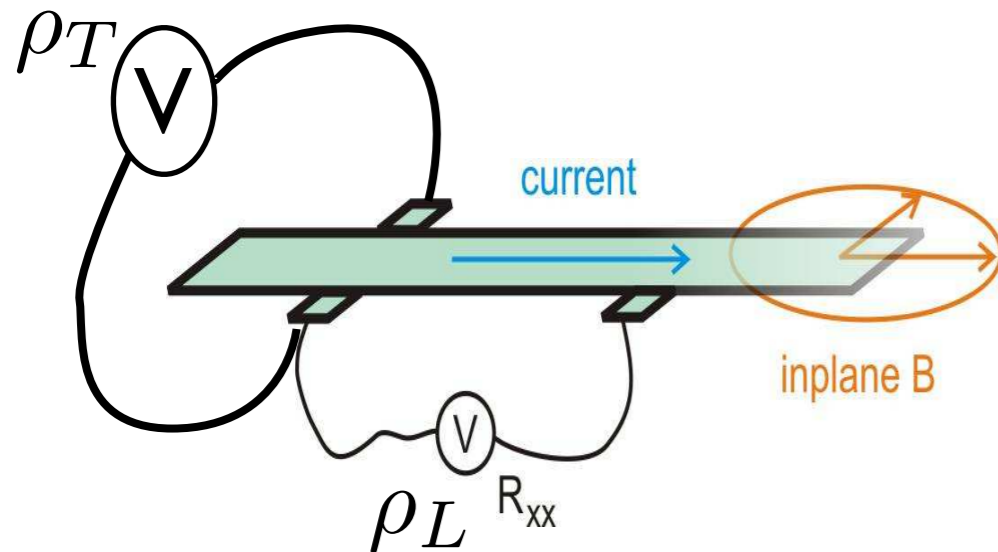
W. Thomson, Proc. Roy. Soc. London (1857)	Fe, Ni	~1%
P. Wiśniewski, Appl. Phys. Lett. (2007)	U <sub>3</sub> As <sub>4</sub>	~50%
A. Rushforth, Phys. Rev. Lett. (2007)	GaMnAs	few %

# AMR: crystalline and non-crystalline components

$$\Delta\rho_L/\rho_{av} = C_I \cos 2\phi + C_{I,C} \cos(2\phi + 4\theta) + C_C \cos(4\phi + 4\theta) + C_U \cos(2\phi + 2\theta)$$

$$\Delta\rho_T/\rho_{av} = C_I \sin 2\phi - C_{I,C} \sin(2\phi + 4\theta)$$

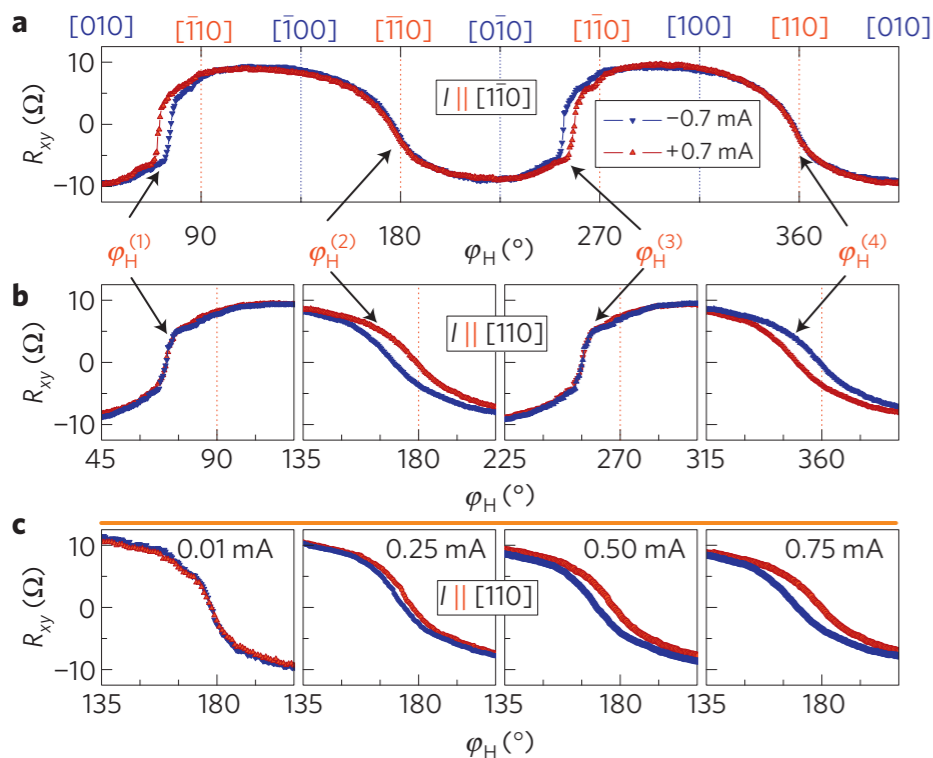
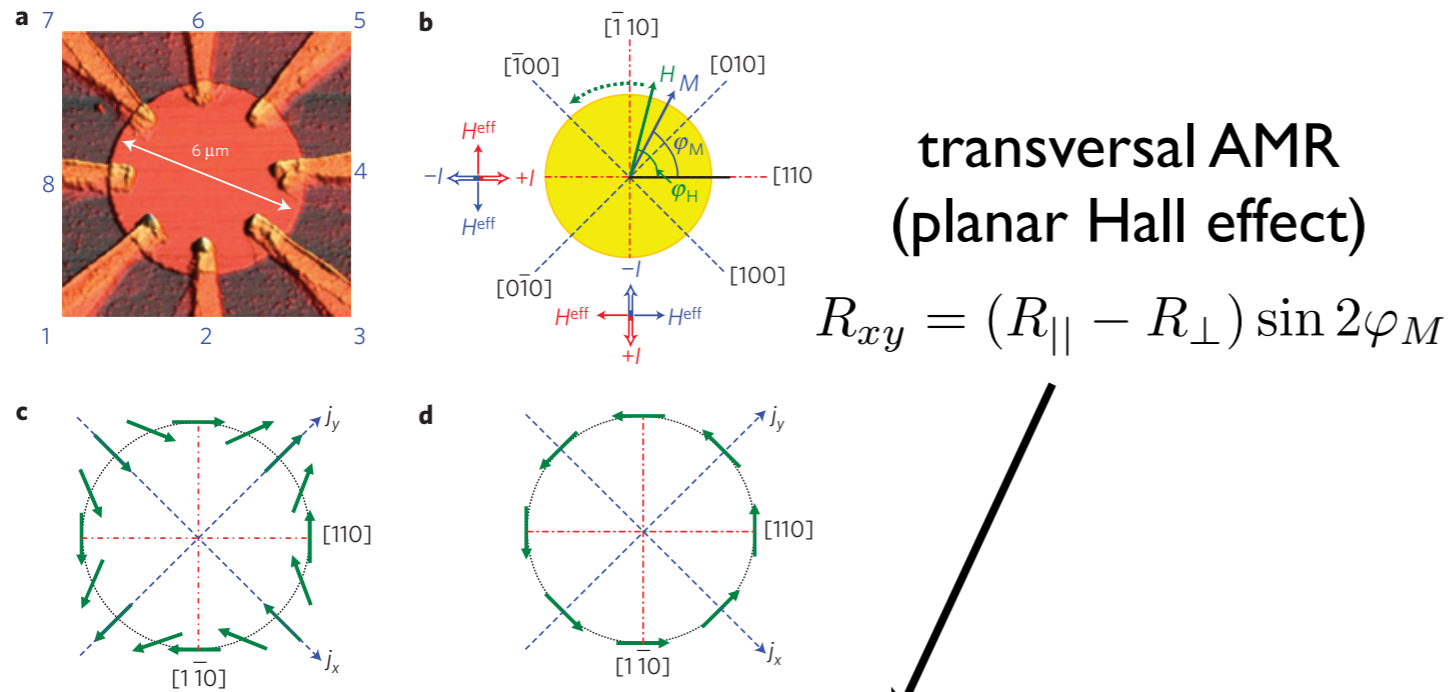
de Ranieri et al., NJP '08





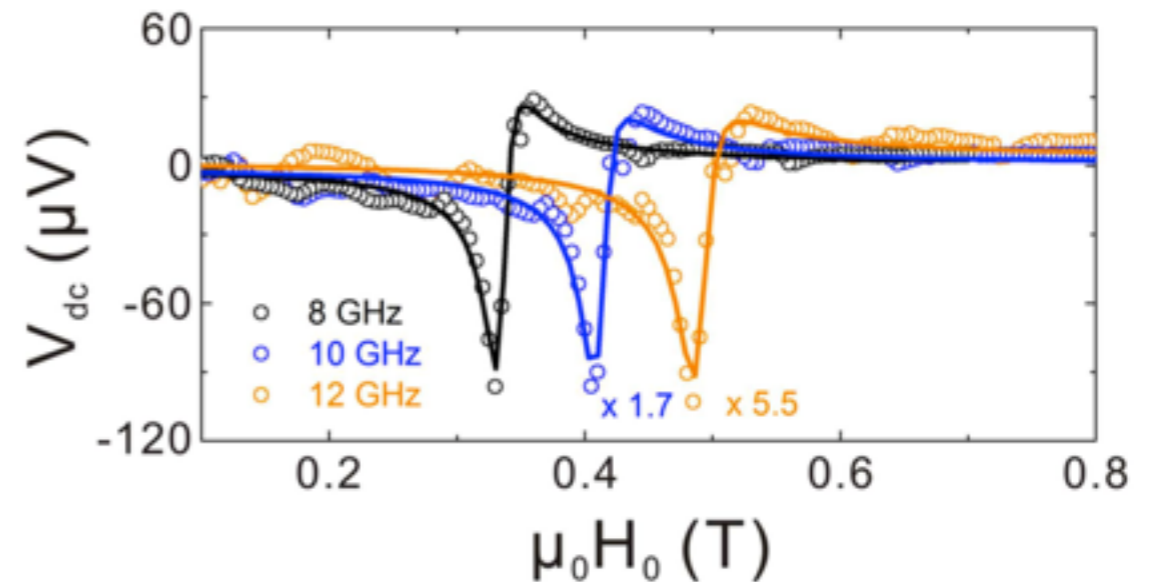
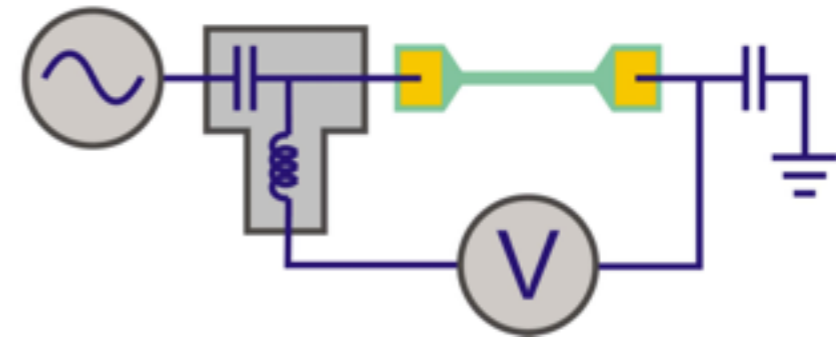
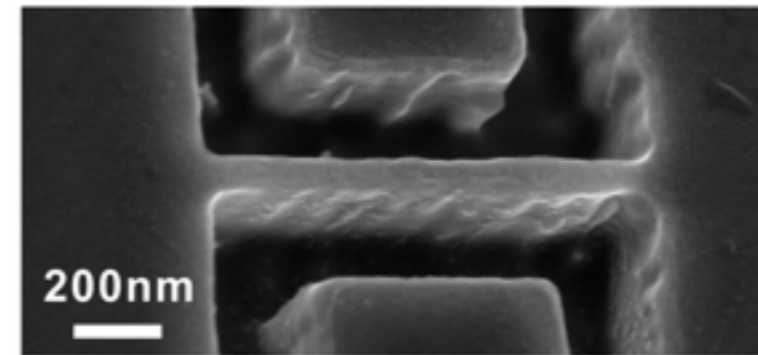
# Current-induced field - experiment in (Ga,Mn)As

static case



Chernyshev et al., Nat Phys (2009)

dynamic case



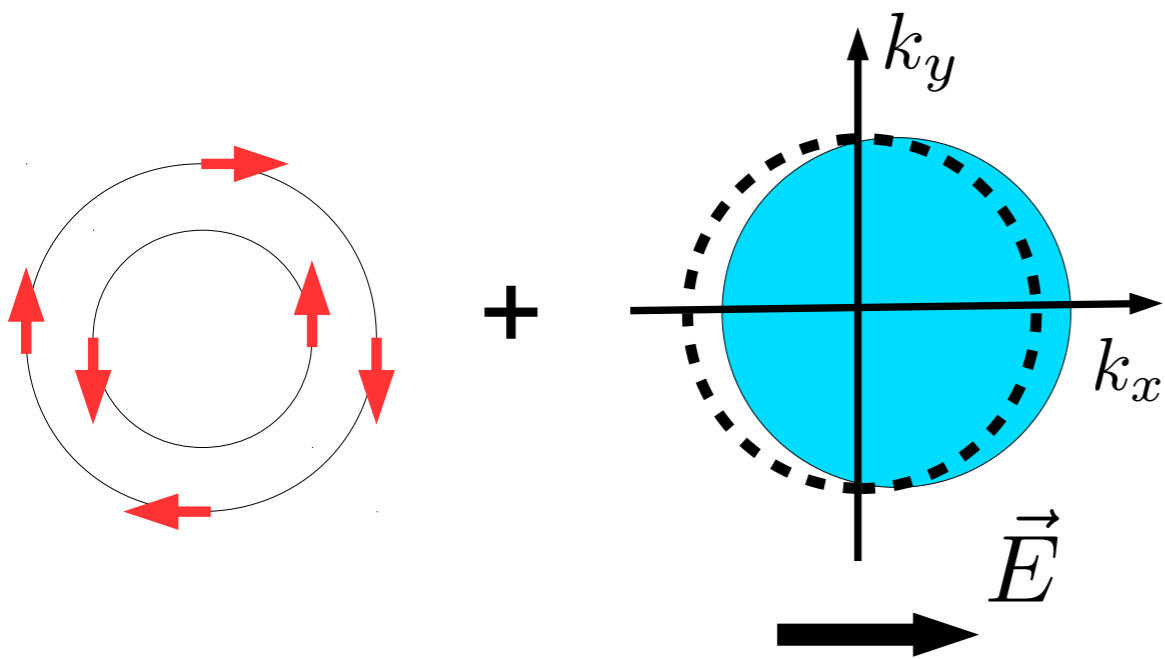
Fang et al., Nat Nano (2011)

# Current-induced spin-orbit torque

Edelstein effect...

$$\delta \mathbf{S} = \chi \mathbf{E}$$

simplest example: Rashba-Bychkov spin-orbit int. (Sol. st. comm. 73, 233)



$$\delta \mathbf{S} = \int S \delta f dk + \int f \delta S dk$$

... action on magnetic moments

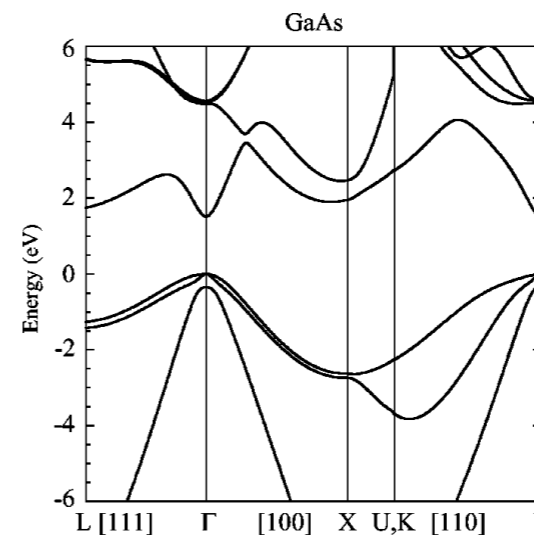
Phys. Rev. B 79, 094422 (2009)

$$\vec{T} = \frac{J_{ex}}{\hbar} \vec{m} \times \vec{M}$$

in the context of s-d type Hamiltonian

$$H = H_{KL} + \hbar \hat{e}_M \cdot \mathbf{s}$$

... applicable to (Ga,Mn)As



- ferromagnetism ind. by carriers
- Mn d-states coupled to hole p-states (carrier)

$$H = H_{KL} + J_{pd} \sum_{i,I} \vec{S}_I \cdot \vec{s}_i \delta(\vec{r}_i - \vec{R}_I)$$

# Current-induced spin-orbit torque in (Ga,Mn)As

## ... a long history

PHYSICAL REVIEW B **80**, 134403 (2009)

### **Influence of a transport current on magnetic anisotropy in gyrotropic ferromagnets**

Ion Garate and A. H. MacDonald

*Department of Physics, The University of Texas at Austin, Austin, Texas 78712, USA*

(Received 23 May 2009; revised manuscript received 5 September 2009; published 5 October 2009)

PHYSICAL REVIEW B **91**, 134402 (2015)

### **Intraband and interband spin-orbit torques in noncentrosymmetric ferromagnets**

Hang Li,<sup>1,\*</sup> H. Gao,<sup>2,3,\*</sup> Liviu P. Zârbo,<sup>4,5</sup> K. Výborný,<sup>5,†</sup> Xuhui Wang,<sup>1</sup> Ion Garate,<sup>6</sup> Fatih Doğan,<sup>1</sup> A. Čejchan,<sup>5</sup>  
Jairo Sinova,<sup>2,3,5</sup> T. Jungwirth,<sup>5,7</sup> and Aurélien Manchon<sup>1,‡</sup>

<sup>1</sup>*King Abdullah University of Science and Technology (KAUST), Physical Science and Engineering Division,  
Thuwal 23955-6900, Saudi Arabia*

<sup>2</sup>*Department of Physics, Texas A&M University, College Station, Texas 77843-4242, USA*

<sup>3</sup>*Institute of Physics, Johannes Gutenberg Universität, 55128 Mainz, Germany*

<sup>4</sup>*Molecular and Biomolecular Physics Department, National Institute for Research and Development of Isotopic  
and Molecular Technologies, RO-400293 Cluj-Napoca, Romania*

<sup>5</sup>*Institute of Physics ASCR, v.v.i., Cukrovarnická 10, 162 53 Praha 6, Czech Republic*

<sup>6</sup>*Département de Physique and Regroupement Québécois sur les Matériaux de Pointe,  
Université de Sherbrooke, Sherbrooke, Québec, Canada J1K 2R1*

<sup>7</sup>*School of Physics and Astronomy, University of Nottingham, Nottingham NG7 2RD, United Kingdom*

(Received 12 January 2015; revised manuscript received 12 March 2015; published 1 April 2015)

# Current-induced spin-orbit torque

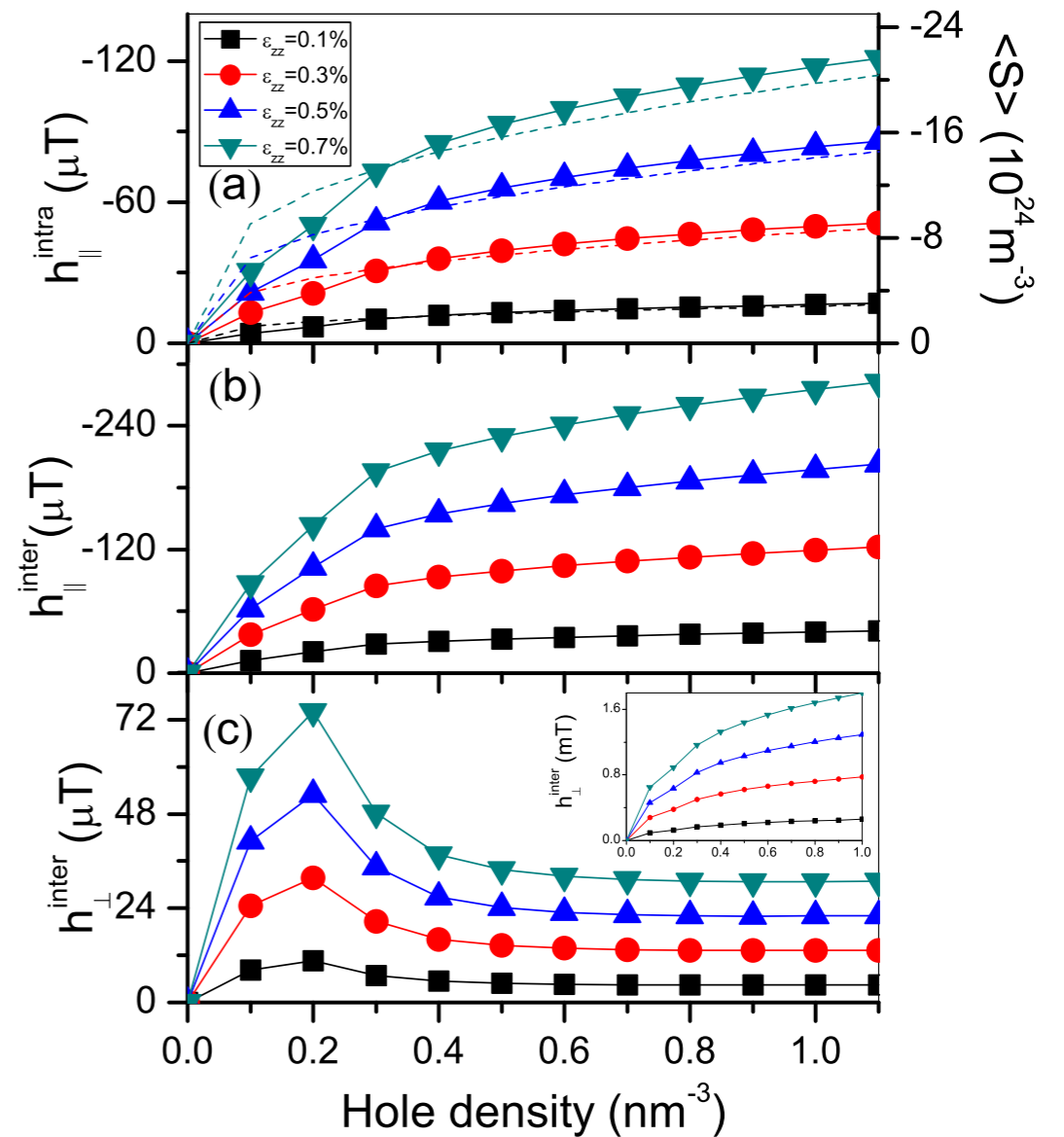
$\delta\mathbf{S} = \chi \mathbf{E}$       Current-induced (non-equilibrium)  
 $S_i = \chi_{ij} E_j$       spin polarisation - in linear response

$$\delta\mathbf{S} = \delta\mathbf{S}^{\text{intra}} + \delta\mathbf{S}_1^{\text{inter}} + \delta\mathbf{S}_2^{\text{inter}}$$

$$\delta\mathbf{S}^{\text{intra}} = \frac{1}{V} \frac{e\hbar}{2\Gamma} \sum_{\mathbf{k},a} \langle \psi_{\mathbf{k}a} | \hat{\mathbf{S}} | \psi_{\mathbf{k}a} \rangle \langle \psi_{\mathbf{k}a} | \mathbf{E} \cdot \hat{\mathbf{v}} | \psi_{\mathbf{k}a} \rangle \times \delta(E_{\mathbf{k}a} - E_F), \quad (3)$$

$$\delta\mathbf{S}_1^{\text{inter}} = -\frac{e\hbar}{V} \sum_{\mathbf{k},a \neq b} 2\text{Re}[\langle \psi_{\mathbf{k}a} | \hat{\mathbf{S}} | \psi_{\mathbf{k}b} \rangle \langle \psi_{\mathbf{k}b} | \mathbf{E} \cdot \hat{\mathbf{v}} | \psi_{\mathbf{k}a} \rangle] \times \frac{\Gamma(E_{\mathbf{k}a} - E_{\mathbf{k}b})}{[(E_{\mathbf{k}a} - E_{\mathbf{k}b})^2 + \Gamma^2]^2} (f_{\mathbf{k}a} - f_{\mathbf{k}b}), \quad (4)$$

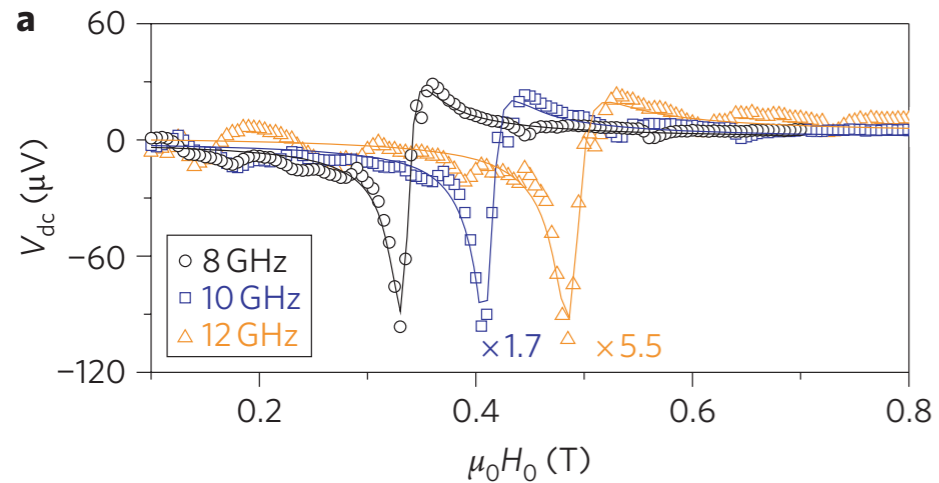
$$\delta\mathbf{S}_2^{\text{inter}} = -\frac{e\hbar}{V} \sum_{\mathbf{k},a \neq b} \text{Im}[\langle \psi_{\mathbf{k}a} | \hat{\mathbf{S}} | \psi_{\mathbf{k}b} \rangle \langle \psi_{\mathbf{k}b} | \mathbf{E} \cdot \hat{\mathbf{v}} | \psi_{\mathbf{k}a} \rangle] \times \frac{\Gamma^2 - (E_{\mathbf{k}a} - E_{\mathbf{k}b})^2}{[(E_{\mathbf{k}a} - E_{\mathbf{k}b})^2 + \Gamma^2]^2} (f_{\mathbf{k}a} - f_{\mathbf{k}b}). \quad (5)$$



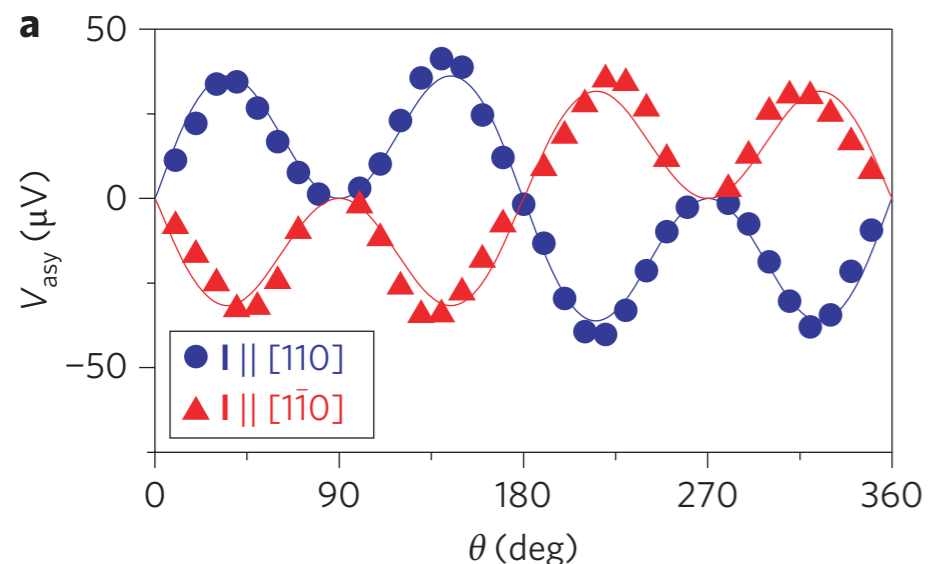
# Current-induced spin-orbit torque

intraband

Nat. Nanotech. 6, 413 (2011)

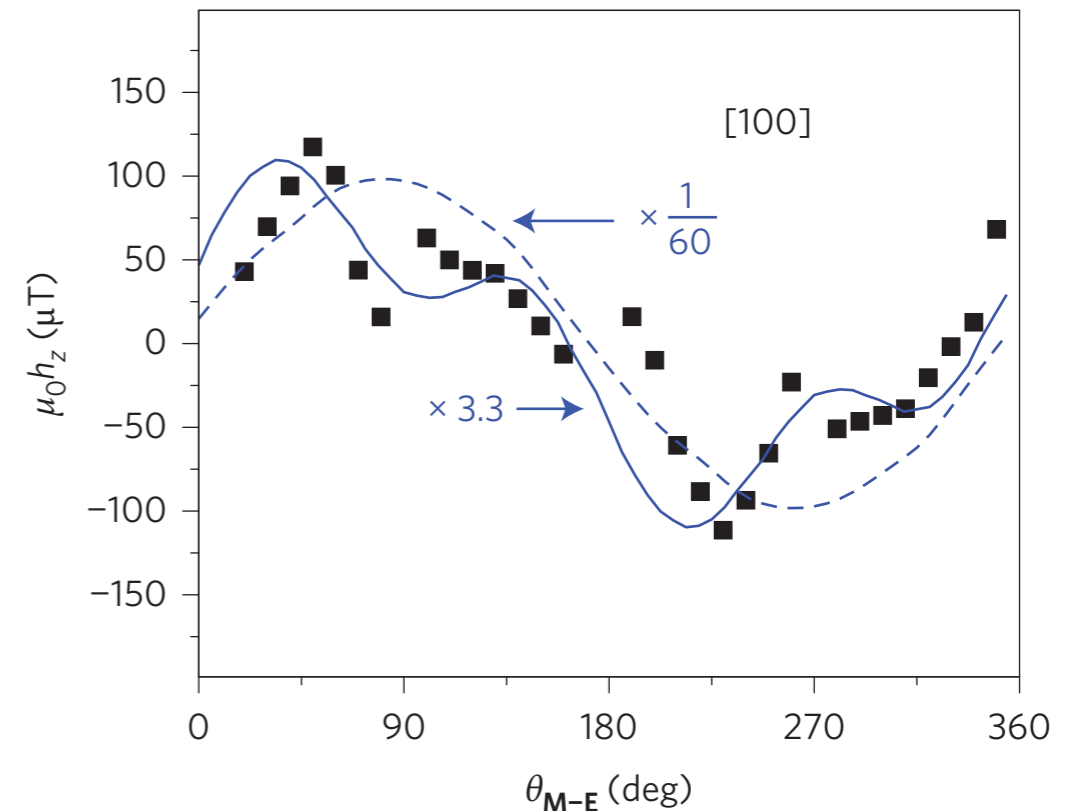


$$\text{Re}\{V_{dc}\} = V_{\text{sym}} \frac{\Delta H^2}{(H_0 - H_{\text{res}})^2 + \Delta H^2} + V_{\text{asy}} \frac{\Delta H(H_0 - \Delta H)}{(H_0 - H_{\text{res}})^2 + \Delta H^2}$$



interband

Nat. Nanotech. 9, 211 (2014)



Berry-curvature type expressions...

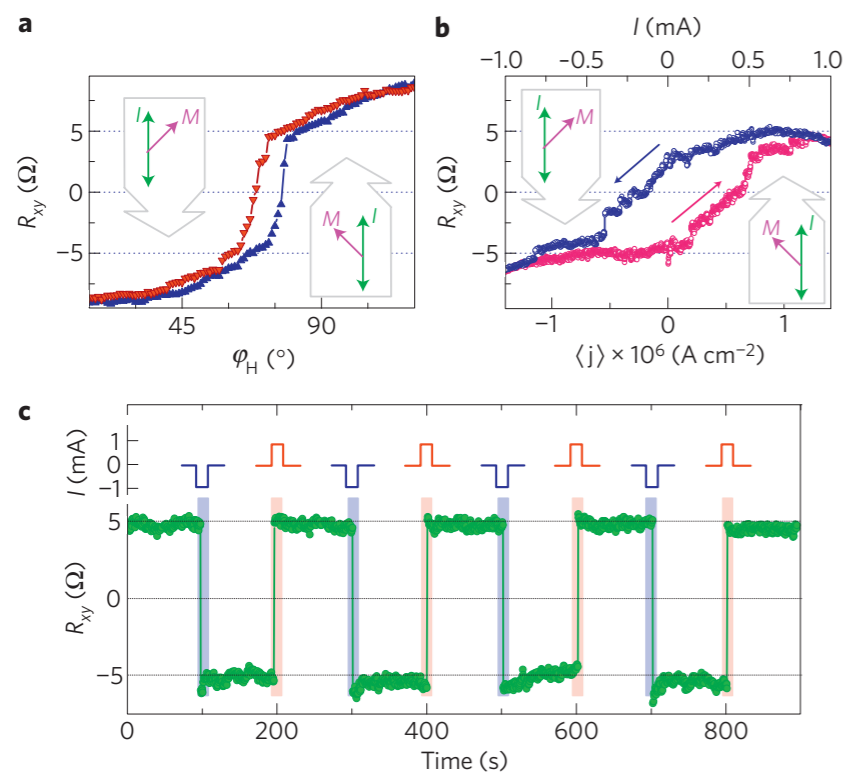
$$S_z = \frac{\hbar}{V} \sum_{\mathbf{k}, a \neq b} (f_{\mathbf{k}, a} - f_{\mathbf{k}, b}) \frac{\text{Im}[\langle \mathbf{k}, a | s_z | \mathbf{k}, b \rangle \langle \mathbf{k}, b | e\mathbf{E} \cdot \mathbf{v} | \mathbf{k}, a \rangle]}{(E_{\mathbf{k}, a} - E_{\mathbf{k}, b})^2}$$

$$S_z = ieE \sum_{\text{occ. states}} \left[ \left\langle \frac{\partial \psi}{\partial h_z} \left| \frac{\partial \psi}{\partial k_x} \right\rangle - \left\langle \frac{\partial \psi}{\partial k_x} \left| \frac{\partial \psi}{\partial h_z} \right\rangle \right]$$

# (Partial) summary

## magnetisation switching by current-induced spin-orbit torque

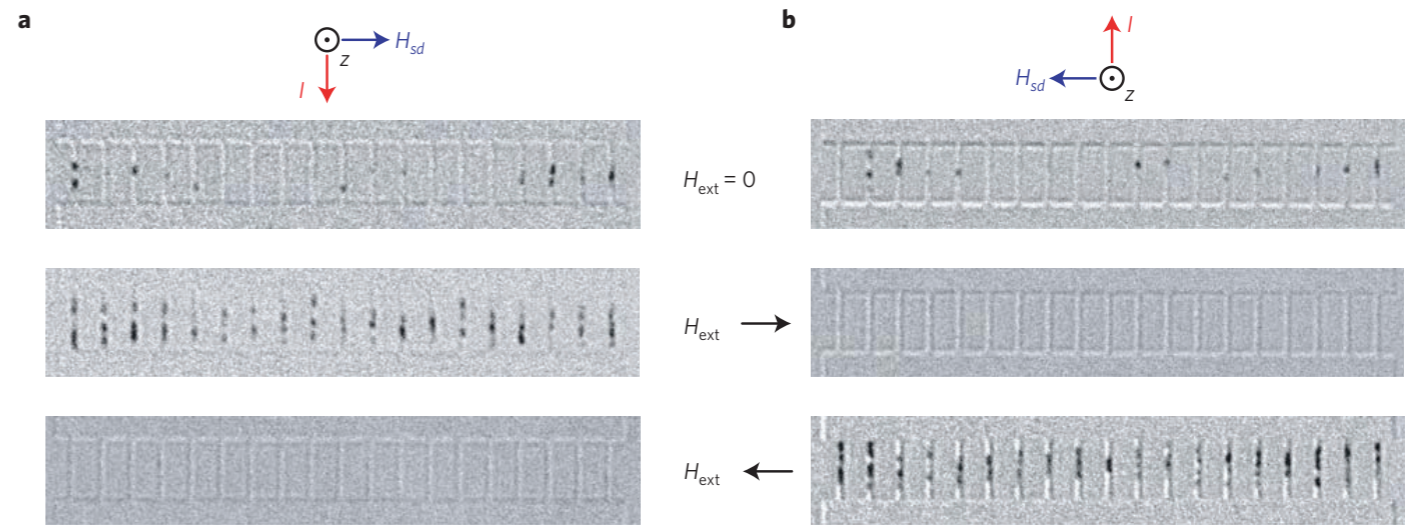
(Ga,Mn)As



**Figure 4 | Current-induced reversible magnetization switching.** **a**,  $\varphi_H$  dependence of  $R_{xy}$  near the  $[010] \rightarrow [\bar{1}00]$  magnetization switching for  $I = \pm 0.7$  mA in sample A for  $I \parallel [\bar{1}10]$ . **b**,  $R_{xy}$  shows hysteresis as a function of current for a fixed field  $H = 6$  mT applied at  $\varphi_H = 72^\circ$ . **c**, Magnetization switches between the  $[010]$  and  $[\bar{1}00]$  directions when alternating  $\pm 1.0$  mA current pulses are applied. The pulses have 100 ms duration and are shown schematically above the data curve.  $R_{xy}$  is measured with  $I = 10 \mu\text{A}$ .

Nat. Phys. 5, 656 (2009)

multilayers - CoPt



**Figure 2 | Differential Kerr microscopy images recorded after current pulse injection.** **a, b**, Positive (**a**) and negative (**b**) current values with external field  $\mu_0 H_{ext} = 0, \pm 47.5$  mT applied parallel to  $\hat{y}$ . The pulse amplitude is set to  $7.8 \times 10^7$  A cm $^{-2}$  in all cases. Note that the topographic contrast in each image varies depending on the drift between pre-pulse and post-pulse images.

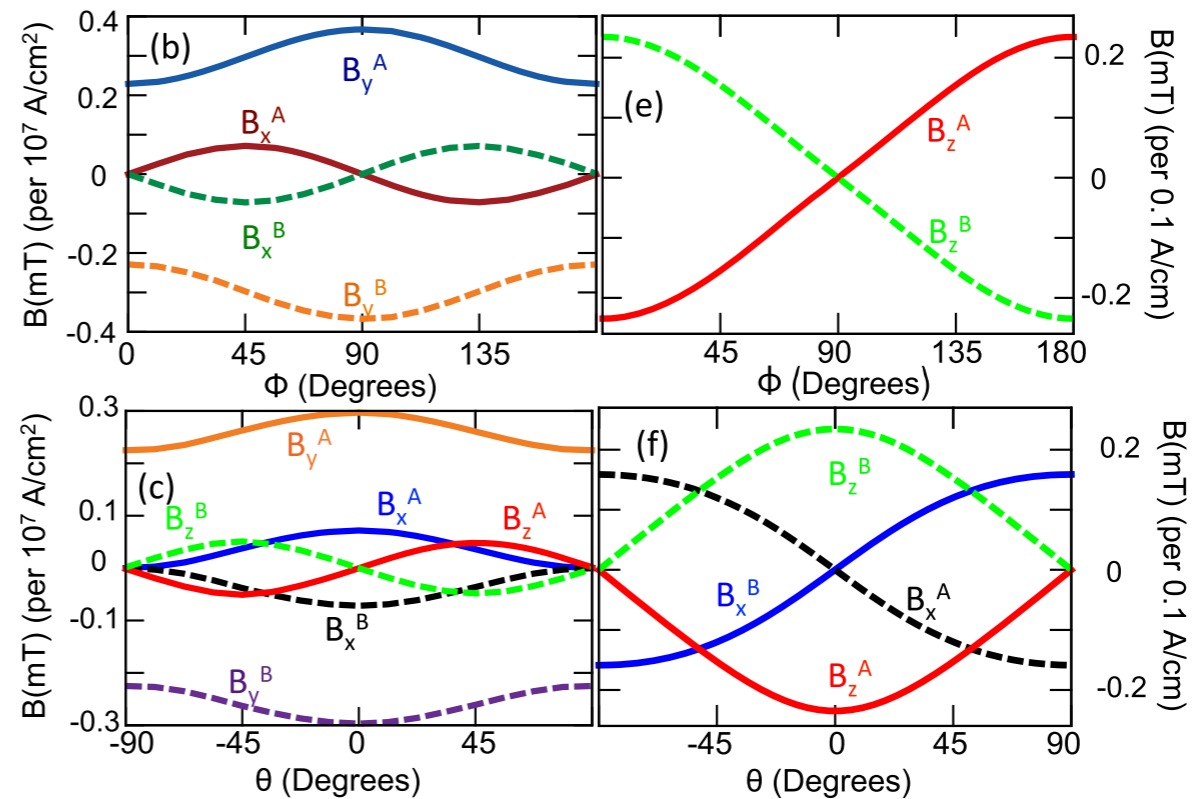
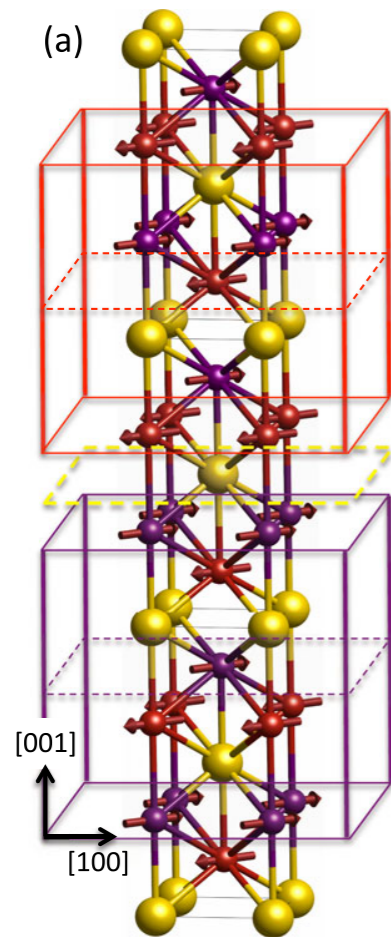
Nat. Mater. 9, 230 (2010)



So far: ferromagnets

Now: antiferromagnets

example:  $\text{Mn}_2\text{Au}$  (Phys. Rev. Lett. 113, 157201)



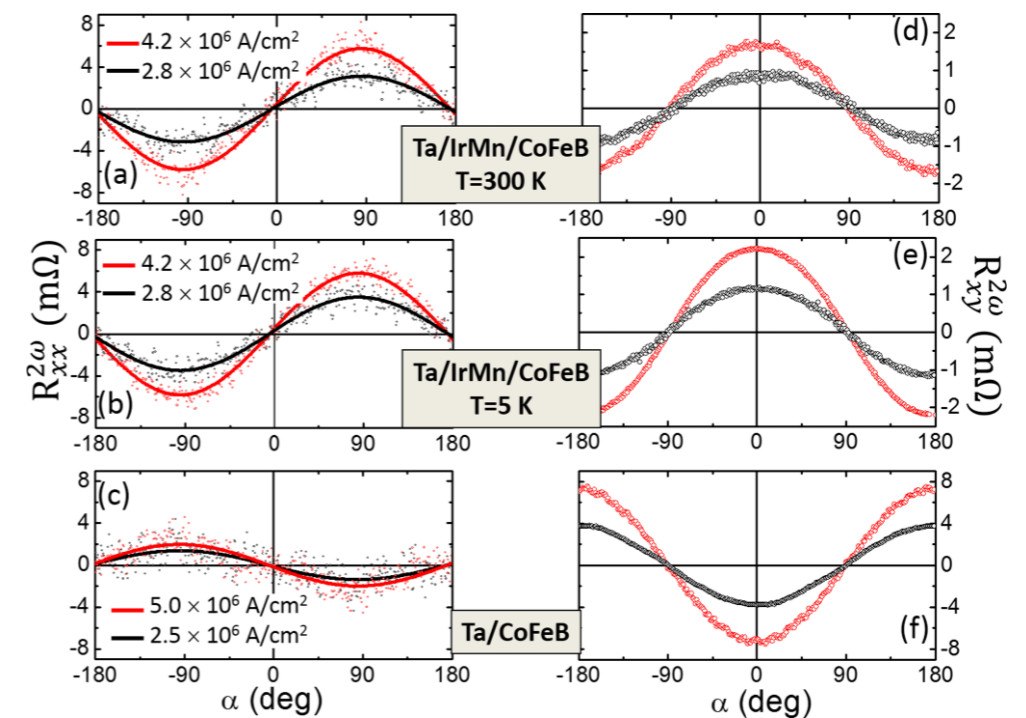
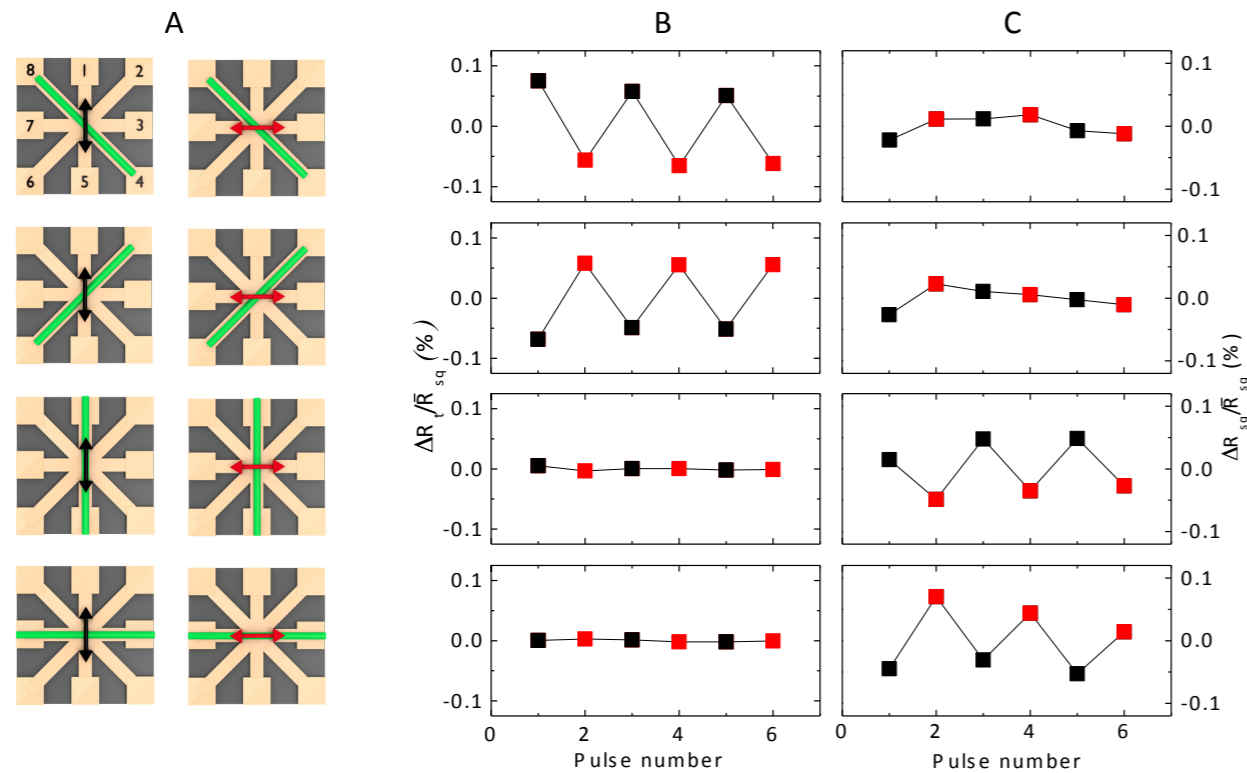
- two sublattices, opposite mag. moments

- current-induced fields resolved by sublattice

# Current-induced spin-orbit torque in antiferromagnets

switching observed  
in CuMnAs

difficult to discriminate  
from other effects (IrMn)



arXiv1503.03765  
(to appear in Science)

Phys. Rev. B 92, 165424 (2015)



# Summary

- current-induced spin-orbit torques observed in various metals
- microscopic theory available
- can be (potentially) used to write information in memory devices
- works both in ferromagnets and antiferromagnets